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How to estimate past quantum measurement interventions after continuous monitoring

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Abstract

We analyze the problem of estimating past quantum states of a monitored system from a mathematical perspective in order to ensure self-consistency with the principle of quantum non-demolition. Despite several claims of “measuring noncommuting observables” in the physics literature, we show that we are always measuring commuting processes. Our main interest is in the notion of quantum smoothing, or retrodiction. In particular, we examine proposals to estimate the result of an external measurement made on an open quantum systems during a period where it is also undergoing continuous monitoring. A full analysis shows that the non-demolition principle is in place, and so a well-posed statistical inference problem can be formulated. We extend the formalism to consider multiple independent external measurements made on the system over the course of a continual period of monitoring.

1 Introduction

Statistical inference uses Bayes Theorem to calculate estimates for unobserved variables in terms of correlated measured ones. However, this may be applied to quantum theory *only* in the case where all these quantities are compatible observables. This is described by the non-demolition principle of V.P. Belavkin [1] which we recall here in the introductory section. This shows that it is permissible to estimate the value of a present (filtering) or a future (prediction) observable of an open system based on past measurement of the output processes as these all commuting. But past observables need not commute with future outputs so their estimation (quantum smoothing) violates the non-demolition principle and constitutes a potential misapplication of Bayes Theorem.

Nevertheless, there have been several proposals for quantum smoothing [3]-[6] addressing the question of whether useful information about the past may still be extracted from future measurements. Here we revisit the proposal by Gammelmark, Julsgaard, and K. Mølmer [5] who consider the situation of an system continuously monitored over the time interval $[0, T]$ upon which we make an external instantaneous measurement at some intermediate time τ . This has

been applied to experimental schemes monitoring superconducting qubits [7]-[11].

The external measurement is taken to be indirect - that is, we follow von Neumann's scheme of preparing a quantum probe in a state, coupling it to the system at time τ and then making an instantaneous measurement of a probe observable. It is assumed that the probe is initially not entangled with either the system or its environment, and that it only couples to the system itself - and not directly to the environment. We refer to such measurements as external *measurement interventions* on the open system. Clearly such interventions perturb the dynamics so that not only is the system affected, but the intervention will also be apparent in the continuous measurement readout. The outcome of the intervention measurement may be concealed, so one may ask for the probability of its outcome conditional on the entire continuously monitored output - both past and future.

The central question is whether such a scheme satisfies the non-demolition principle. If it does not, then we have the situation that the measurements made beyond time τ degrade the estimate since we would now be measuring incompatible observables, and theoretically we have no ground to be applying Bayes Theorem. Following a careful examination of what commutes with what, our main result is to show that the scheme does actually satisfy the non-demolition principle and to use standard quantum filtering theory to provided the least squares estimate for the probe's observable. This re-derives the results of Gammelmark, *et al.* [5]. In section 3, we extend the situation to several independent measurement interventions made over the course of the time interval $[0, T]$.

1.1 Instantaneous Measurements

A simple description of indirect measurement is the following. Suppose our system has Hilbert space \mathfrak{h} and is prepared in a state with density matrix ρ . We set up a probe with Hilbert space \mathfrak{H} and which has been independently prepared in a state ρ_{probe} , then we couple it to the system using some unitary V . We measure an observable M of the probe which we take to have spectral decomposition $M = \sum_m P_m$. If we measure the eigenvalue m for the probe observable M then the state of the system updates to

$$\rho \mapsto \frac{1}{\text{tr}\{\Phi_m(\rho)\}} \Phi_m(\rho). \quad (1)$$

where Φ_m is the CP map

$$\Phi_m(\rho) = \text{tr}_{\mathfrak{H}}\{V\rho \otimes \rho_{\text{probe}}V^* I \otimes P_m\}. \quad (2)$$

As a special case, let $\{|m\rangle\}$ form an orthonormal basis for the probe Hilbert space and take the probe to be in the pure state $|\varphi\rangle$. Suppose that, for each $|\psi\rangle$ in \mathfrak{h} , we have $V|\psi\rangle \otimes |\varphi\rangle = \sum_m V_m|\psi\rangle \otimes |m\rangle$ where the V_m 's are operators on the system space. Then $\Phi_m(\rho) \equiv V_m\rho V_m^*$.

$$\rho \mapsto \frac{1}{\text{tr}\{\rho V_m^* V_m\}} V_m \rho V_m^* \quad (3)$$

There are two simple rules that need to be followed when trying to extract quantum information from measurements. Rule # 1 (Self-Non-Demolition Principle): All the observables measured should be compatible with each other. Rule # 2 (Non-Demolition Principle): All the observables estimated should be compatible with all measured observables.

If we wish to infer useful information about an observable X from measured observables $\{Y_1, \dots, Y_N\}$ and, then the principles combined will require that the family $\{X, Y_1, \dots, Y_N\}$ is commutative. In this case there exist well-defined joint probabilities and Bayes Theorem may be applied to condition X on the observations. If the rules are not followed then we are misapplying Bayes Theorem.

In general, given a state $\mathbb{E} = \text{tr}\{\rho\}$ we may define the conditional expectation of X given Y provided they are compatible. Let $P_X[dx]$ and $P_Y[dy]$ be the projection-valued measures associate with X and Y respectively. In this case $p_{X,Y}[dx, dy] = \text{tr}\{\rho P_X[dx] P_Y[dy]\}$ is a well-defined joint probability measure for both observables and the conditional probability $p_{X|Y}[dx|y]$ is well-defined via $p_{X|Y}[dx|y] p_Y[dy] \equiv p_{X,Y}[dx, dy]$ we set

$$\mathbb{E}[X|Y] = \int x \hat{P}_X[dx] \quad (4)$$

where

$$\hat{P}_X[dx] = \int p_{X|Y}[dx|y] P_Y[dy]. \quad (5)$$

In principle, for a given state \mathbb{E} , we can always define the conditional expectation $\mathbb{E}[X|\mathfrak{Y}]$ of an observables X onto the algebra \mathfrak{Y} generated by the measured observables provided that the non-demolition rules above hold.

1.2 Quantum Filtering

Let us consider a fairly standard set up in quantum trajectories. We have a system with state space \mathfrak{h} coupled to a bath \mathfrak{F} . The bath consists of n input processes $b_{\text{in},k}(t)$ satisfying singular commutation relations $[b_{\text{in},j}(t), b_{\text{in},k}(s)^*] = \delta_{jk} \delta(t-s)$ and we assume the vacuum state $|\Omega\rangle$ for the bath. The coupling of the bath to the system over the time interval t_1 to t_2 is taken to be described by the unitary on $\mathfrak{h} \otimes \mathfrak{F}$

$$U(t_2, t_1) = \vec{T} e^{-i \int_{t_1}^{t_2} \Upsilon(s) ds} \quad (6)$$

where $-i\Upsilon(t) = \sum_k \{L_k \otimes b_{\text{in},k}(t)^* - L_k^* \otimes b_{\text{in},k}(t)\} - iH$. We note the flow property

$$U(t_3, t_2) U(t_2, t_1) = U(t_3, t_1), \quad (t_3 \geq t_2 \geq t_1). \quad (7)$$

We intend to measure the field quadratures. To this end, set $Z_k(t) = \int_0^t \{b_{\text{in},k}(s) + b_{\text{in},k}(s)^*\} ds$, then we note that $[Z_j(t), Z_k(s)] = 0$. However, what we must measure are the outputs

$$Y_k(t) = U(t, 0)^* \{I \otimes Z_k(t)\} U(t, 0) = \int_0^t \{b_{\text{out},k}(s) + b_{\text{out},k}(s)^*\} ds. \quad (8)$$

One may show that $b_{\text{out},k}(t) = I \otimes b_{\text{in},k}(t) + j_t(L_k)$ where, for any system operator X ,

$$j_t(X) = U(t, 0)^* \{X \otimes I\} U(t, 0). \quad (9)$$

We now recall some basic facts [1].

Proposition 1 *The family $\{Y_k(t) : k = 1, \dots, n, 0 \leq t \leq T\}$ is a commutative family. The non-demolition principle holds for $j_t(X)$ and the measured observables $\{Y_k(u) : k = 1, \dots, n, 0 \leq u \leq s\}$ when $t \geq s$.*

To see this, suppose that $t \geq s$, then we observe that

$$Y_k(s) = U(t, 0)^* \{I \otimes Z_k(s)\} U(t, 0). \quad (10)$$

as the right hand side is $U(s, 0)^* U(t, s)^* \{I \otimes Z_k(s)\} U(t, s) U(s, 0)$, however $U(t, s)$ couples the system to only those input processes over the time interval s to t and so commutes with $I \otimes Z_k(t)$, and by unitarity $U(t, s)^* U(t, s)$ is the identity. We therefore have that

$$[Y_j(t), Y_k(s)] = U(t, 0)^* \{I \otimes [Z_j(t), Z_k(s)]\} U(t, 0) \quad (11)$$

which vanishes as the Z operators commutes. Likewise, for $t \geq s$ we may use (10) again to show that

$$[j_t(X), Y_j(s)] = U(t, 0)^* [X \otimes I, I \otimes Z_k(s)] U(t, 0) = 0.$$

We may therefore estimate the present (filtered), or future (predicted), value of a system observable based on the quadrature observations up to the present time. Estimating past values (smoothing) is not possible as it violates the non-demolition principle.

The filter is well-known for this problem. The best estimate for $j_t(X)$ is the conditional expectation

$$\pi_t(X) = \mathbb{E}[j_t(X) | \mathfrak{Y}_{[0,t]}], \quad (12)$$

where $\mathfrak{Y}_{[0,t]}$ is the algebra generated by $\{Y_k(u) : k = 1, \dots, n, 0 \leq u \leq t\}$.

We may write $\pi_t(X) = \text{tr}\{\hat{\rho}_t X\}$ where $\hat{\rho}_t$ satisfies the stochastic partial differential equation (SPDE)

$$d\hat{\rho}_t = i[H, \hat{\rho}_t] dt + \mathcal{D}(\hat{\rho}_t) dt + \sum_k (L_k \hat{\rho}_t + \hat{\rho}_t L_k^* - \lambda_k(t) \hat{\rho}_k) dI_k(t), \quad (13)$$

with $\mathcal{D}(\hat{\rho}_t) = \sum_k (L_k \hat{\rho}_t L_k^* - \frac{1}{2} L_k^* L_k \hat{\rho}_t - \frac{1}{2} \hat{\rho}_t L_k^* L_k)$, $\lambda_k(t) = \text{tr} \{ \hat{\rho}_t (L_k + L_k^*) \}$ and the $I_k(t)$ are Wiener processes known as the *innovations processes* and are given by

$$dI_k(t) = dY_k(t) - \sqrt{\eta_k} \lambda_k(t) dt, \quad (14)$$

where we allow for an efficiency factor $\eta_k \in [0, 1]$ for the k th measurement. The derivation of the filter makes explicit use of Bayes Theorem.

We remark that the SPDE (13) is nonlinear in $\hat{\rho}_t$, however, it is advantageous to work with an unnormalized version which obeys a linear SDPE. It is possible to write $\hat{\rho}_t = \rho_t / \text{tr} \{ \rho_t \}$ where ρ_t satisfies

$$d\rho_t = i[H, \rho_t] dt + \mathcal{D}(\rho_t) dt + \sum_k (L_k \rho_t + \rho_t L_k^*) dY_k(t) \quad (15)$$

Remark 2 The filter $\hat{\rho}_t$ is to be obtained by solving the SPDE (13) subject to an initial condition $\hat{\rho}_0 = \rho_0$. The initial state ρ_0 may not always be known in which case guessing at one may lead to error. In many cases, the sensitivity to initial conditions is not important and a wrongly initialized filter will converge to the correct value asymptotically in time. See, for instance, [12, 13, 14].

Remark 3 From the readout Y_k we may extract $\lambda_k(t) = \text{tr} \{ \hat{\rho}_t (L_k + L_k^*) \}$. This is colloquially phrased as “measuring the observables $L_k + L_k^*$ ”. As the $L_k + L_k^*$ need not commute, it gives the impression that one is measuring noncommuting observables. However, what happens is that the field quadratures, $Y_k(t)$, are measured and these do commute! One then estimates the expected value of $L_k + L_k^*$ from this. In reality, Proposition 1 shows that the quadratures $Y_k(t)$ commute for all k and all $t \geq 0$. So the totality of what we measure is compatible - however, we can use this to estimate noncommuting observables.

1.3 Computing Estimates

Let us now sketch the argument for deriving the estimates. We adapt the presentation given in [15].

Theorem 4 (Bouten-van Handel) Let $\mathbb{E} = \text{tr} \{ \rho \}$ be a state on an algebra \mathfrak{A} of operators and \mathfrak{Y} a commutative subalgebra generated by the measured observables in a given experiment. Suppose that $\mathfrak{Y} = U^* \mathfrak{Z} U$ for a given unitary U . Furthermore, suppose that there is a fixed $F \in \mathfrak{Y}'$ such that, for every $X \in \mathfrak{Y}'$, we have $\mathbb{E}[U^* X U] = \mathbb{E}[F^* X F]$. Then

$$\mathbb{E}[U^* X U | \mathfrak{Y}] = \frac{1}{\sigma(I)} \sigma(X), \quad (16)$$

where $\sigma(X) = U^* \mathbb{E}[F^* X F] \mathfrak{Z} U$.

To see how to apply this to the filtering problem, let us first of all note that the algebra $\mathfrak{Y}_{[0,t]}$ generated by the output quadratures over the time interval

$[0, t]$ may be written as $U(t)^*[I \otimes \mathfrak{Z}_{[0,t]}]U(t)$ by virtue of (10) from Proposition 1. We next observe that we may write

$$\mathbb{E}[U(t, 0)^*[X \otimes I]U(t, 0)] = \mathbb{E}[F(t, 0)^*[X \otimes I]F(t, 0)]$$

where we introduce the two-parameter family of time-ordered exponentials

$$F(t_2, t_1) = \vec{T} e^{\int_{t_1}^{t_2} [\sum_k L_k dZ_k(t) + K dt]}, \quad (17)$$

where $K = -\frac{1}{2} \sum_k L_k^* L_k - iH$. Here the state is a tensor product of system and Fock vacuum state of the bath: $\rho_{\text{sys}} \otimes |\Omega\rangle\langle\Omega|$. Indeed it is easy to show that

$$U(t_2, t_1) |\psi\rangle \otimes |\Omega\rangle = F(t_2, t_1) |\psi\rangle \otimes |\Omega\rangle,$$

for any system vector $|\psi\rangle$.

We use the Theorem to write $\pi_t(X) = \mathbb{E}[j_t(X)|\mathfrak{Y}_{[0,t]}] = \sigma_t(I)^{-1} \sigma_t(X)$, with $\sigma_t(X) = \text{tr}\{\rho_{\text{sys}} F^Y(t, 0)^* X F^Y(t, 0)\}$, and where we introduce the new process

$$F^Y(t_2, t_1) = \vec{T} e^{\int_{t_1}^{t_2} [\sum_k L_k dY_k(t) + K dt]}, \quad (18)$$

It is not difficult to see that $\sigma_t(X)$ satisfies the stochastic differential equation

$$d\sigma_t(X) = \sigma_t(\mathcal{L}X) dt + \sum_k \sigma_t(X L_k + L_k^* X) dY_k(t). \quad (19)$$

where the (Heisenberg picture) Lindbladian is $\mathcal{L}X = \sum_k L_k^* X L_k + X K + K^* X$.

This is equivalent to (13) for the unit efficiency case. In the classical world, we encounter the Zakai equation which is a *linear* SPDE for the unnormalized filter which occurs in continuous time estimation problems. The equation (19) is its quantum analogue, and we refer to it as the *Belavkin-Zakai equation*.

2 Intervention Measurements

The proposal was made by Gammelmark *et alia* that, in addition to the quadrature measurements over the time interval $[0, T]$, one might also make an intervention measurement at an intermediate time τ ($0 < \tau < T$). The question we ask here is whether we violate the non-demolition principle in the process.

If, rather than an instantaneous measurement at time τ , we are prepared to settle for another indirect measurement then there is a simple way to do this. The additional measurement is realized as a quadrature measurement of an extra input field and so we just have the old filtering problem with a larger number of channels.

However, we want to make an instantaneous measurement at time τ . It is clear that if this is a direct measurement of some observable A of the system then what in practice we are doing is measuring $j_\tau(A)$ which we know does not

commute with the output quadratures beyond time τ . This necessitates using indirect measurements. To this end we introduce an auxiliary system with Hilbert space \mathfrak{H} which will act as the probe. For simplicity, we assume that the auxiliary system is in an initial state ρ_{probe} and has no internal evolution. We now work on the tensor product $\mathfrak{h} \otimes \mathfrak{F} \otimes \mathfrak{H}$, and have the same coupling between the system and bath as before, no coupling between that bath and probe, and no interaction between the system and probe save for a unitary V which is applied at time τ . We understand that $[V, I \otimes F \otimes I] = 0$ for all bath operators F , however, V must lead to a nontrivial coupling between the system and probe so that we may obtain any information about the system by measuring the probe.

The unitary dynamics on $\mathfrak{h} \otimes \mathfrak{F} \otimes \mathfrak{H}$ is now described by

$$\tilde{U}(t, 0) = \begin{cases} U(t, 0) \otimes I, & t < \tau; \\ [U(t, \tau) \otimes I] V [U(\tau, 0) \otimes I], & t \geq \tau. \end{cases} \quad (20)$$

We also set

$$\tilde{U}(t, s) = \tilde{U}(t, 0) \tilde{U}(s, 0)^* \quad (t \geq s).$$

(The two-parameter family $\tilde{U}(t, s)$ therefore satisfies the flow property.)

The output quadratures are now

$$\tilde{Y}_k(t) = \tilde{U}(t, 0)^* [I \otimes Z_k(t) \otimes I] \tilde{U}(t, 0). \quad (21)$$

Likewise, the observable of the probe which is measured will be of the form

$$\begin{aligned} \tilde{M}(\tau) &= \tilde{U}(\tau, 0)^* [I \otimes I \otimes M] \tilde{U}(\tau, 0) \\ &= [U(\tau, 0) \otimes I]^* V^* [I \otimes I \otimes M] V [U(\tau, 0) \otimes I], \end{aligned} \quad (22)$$

for some fixed self-adjoint operator M on \mathfrak{H} . Finally, system observables will have the modified evolution

$$\tilde{J}_t(X) = \tilde{U}(t, 0)^* [X \otimes I \otimes I] \tilde{U}(t, 0). \quad (23)$$

Proposition 5 *The family $\{\tilde{Y}_k(t) : k = 1, \dots, n, 0 \leq t \leq T\}$ is a commutative family.*

Proof. Suppose that $t \geq s$, then similar to (10) we observe that

$$\tilde{Y}_k(s) = \tilde{U}(t, 0)^* \{I \otimes Z_k(s) \otimes I\} \tilde{U}(t, 0). \quad (24)$$

The right hand side is $\tilde{U}(t, 0)^* \tilde{U}(t, s)^* \{I \otimes Z_k(s) \otimes I\} \tilde{U}(t, s) \tilde{U}(t, 0)$.

For $t \geq s \geq \tau$, the unitary $U(t, s)$ couples only the input processes over the time interval s to t to the system and so commutes with $I \otimes Z_k(t) \otimes I$. If $t \leq \tau \leq s$, then we have an extra factor V to worry about, but this only couples the system and probe and so again commutes with $I \otimes Z_k(s) \otimes I$. Finally, if

$\tau > t \geq s$, then we are back in the situation of Proposition 1. We therefore have that $[\widetilde{Y}_j(t), \widetilde{Y}_k(s)]$ equals

$$\widetilde{U}(t, 0)^* \{I \otimes [Z_j(t), Z_k(s)] \otimes I\} \widetilde{U}(t, 0) = 0.$$

■

Proposition 6 *The observable $\widetilde{M}(\tau)$ commutes with the quadrature observables $\{\widetilde{Y}_k(t) : k = 1, \dots, n, 0 \leq t \leq T\}$.*

Proof. We separate into two cases. First, for $t \leq \tau$ we use (24) to write

$$\widetilde{Y}_k(t) = \widetilde{U}(\tau, 0)^* \{I \otimes Z_k(t) \otimes I\} \widetilde{U}(\tau, 0), \quad (25)$$

which obviously commutes with $\widetilde{M}(\tau)$ as given by (22).

Next, for $t \geq \tau$, let us first note that $\widetilde{U}(t, \tau) = U(t, \tau) \otimes I$ couples the system to the bath processes over the interval $[\tau, t]$ and therefore commutes with $I \otimes I \otimes M$ giving

$$\begin{aligned} \widetilde{M}(\tau) &= \widetilde{U}(\tau, 0)^* \widetilde{U}(t, \tau)^* [I \otimes I \otimes M] \widetilde{U}(t, \tau) \widetilde{U}(\tau, 0) \\ &= \widetilde{U}(t, 0)^* [I \otimes I \otimes M] \widetilde{U}(t, 0). \end{aligned} \quad (26)$$

We then see that $[\widetilde{M}(\tau), \widetilde{Y}_j(t)]$ equals

$$\widetilde{U}(t, 0)^* [I \otimes I \otimes M, I \otimes Z_k(t) \otimes I] \widetilde{U}(t, 0) = 0.$$

■

We therefore have that *all* the measurements are compatible, and so establish self-nondemolition. The next question then is what we can now hope to estimate with them.

Proposition 7 *The observable $\widetilde{J}_t(X)$ commutes with the set of quadrature operators $\{\widetilde{Y}_k(u) : k = 1, \dots, n, 0 \leq u \leq t\}$ and the observable $\widetilde{M}(\tau)$ provided $t \geq \tau$.*

Proof. The commutativity of $\widetilde{J}_t(X)$ with the quadratures up to and including time t follows from an argument virtually identical to that in Proposition 1. Next, for $t \geq \tau$, we use (26) to write $[\widetilde{J}_t(X), \widetilde{M}(\tau)]$ as

$$\widetilde{U}(t, 0)^* [X \otimes I \otimes I, I \otimes I \otimes M] \widetilde{U}(t, 0) = 0.$$

■

Outside of this, there are no natural constraints forcing any of the various observables under consideration to commute.

2.1 Summary

The continuously monitored quadratures $\{\tilde{Y}_k(t) : k = 1, \dots, n, t \in [t_1, t_2]\}$ generate a commutative (von Neumann) algebra for all $t_1 < t_2$ which we denote as $\tilde{\mathfrak{Y}}_{[t_1, t_2]}$. This is the *algebra of the quadrature observations* over the time interval $[t_1, t_2]$. We have the isotonic condition

$$\tilde{\mathfrak{Y}}_{[t_1, t_2]} \subset \tilde{\mathfrak{Y}}_{[t_3, t_4]} \quad (27)$$

whenever the interval $[t_1, t_2]$ is contained inside $[t_3, t_4]$. The isotony condition implies that $\{\tilde{\mathfrak{Y}}_{[0, t]} : 0 \leq t \leq T\}$ is a nested family of commutative algebras, known as a *filtration*.

The algebra generated by $\tilde{\mathfrak{Y}}_{[0, T]}$ and the additional observable $\tilde{M}(\tau)$ is again a commutative algebra and this is what we intend to condition into.

To find the best estimate of an observable X of the system at time $t \in [0, T]$, i.e., estimate $\tilde{\pi}_t(X)$, we use the observations of the quadratures up to and including time t and additionally $\tilde{M}(\tau)$ if $t \geq \tau$.

The best estimate will have the form $\tilde{\pi}_t(X) = \text{tr}\{\tilde{\rho}_t X\}$ where

$$\tilde{\rho}_t = \hat{\rho}_t, \quad (t < \tau) \quad (28)$$

where $\hat{\rho}_t$ is the solution to the basic filter SPDE (13) with initial condition ρ_0 at time 0;

$$\tilde{\rho}_\tau = \frac{1}{\text{tr}\{\Phi_m(\hat{\rho}_\tau)\}} \Phi_m(\hat{\rho}_\tau); \quad (29)$$

and, for times $t \geq \tau$, $\tilde{\rho}_t$ will be the solution to the basic filter SPDE (13) with initial condition $\tilde{\rho}_\tau$ as computed in (29) now initialized at time τ .

Remark 8 *At this stage it is fairly obvious that we may extend the theory to multiple intervention measures at times $\tau_1 < \tau_2 < \dots < \tau_r$ during the time interval $[0, T]$. The procedure is simple enough: we use fresh probe systems at each intervention to ensure that all measurement are compatible; we use the basic filter SPDE (13) to propagate the filter starting with state ρ_0 at time 0, then update the state by*

$$\tilde{\rho}_{\tau_k^-} \mapsto \tilde{\rho}_{\tau_k^+} = \frac{1}{\text{tr}\{\Phi_{(m_k)}^{(k)}(\tilde{\rho}_{\tau_k^-})\}} \Phi_{(m_k)}^{(k)}(\tilde{\rho}_{\tau_k^-}) \quad (30)$$

at each time τ_k where m_k is the recorded value at the k measurement (described by the CP maps $\Phi_{(m_k)}^{(k)}$).

2.2 Estimating what the probe measures

In their paper, Gammelmark et al. [5] consider the situation where the result of the probe measurement at time τ is not used - locked in a safe deposit box

until after the quadrature measurement period $[0, T]$ is over. They then ask as for the probability that the intervention measurement was value m conditional on the whole history of the monitored quadratures.

In our language, we would say that they are estimating the observable $\widetilde{M}(\tau)$ using the quadrature measurements, or equivalently conditioning the observable $\widetilde{M}(\tau)$ onto $\mathfrak{Y}_{[0, T]}$.

By Proposition 6, this is possible! Let us take the spectral decomposition $M = \sum_m m P_m$, then the goal is to compute

$$p(m, \tau) = \mathbb{E}[\tilde{U}(\tau, 0)^* [I \otimes I \otimes P_m] \tilde{U}(\tau, 0) | \mathfrak{Y}_{[0, T]}]. \quad (31)$$

This may be rewritten as

$$\tilde{U}(\tau, 0) \mathbb{E}_{[0, T]} [I \otimes I \otimes P_m | I \otimes \tilde{\mathfrak{Z}}_{[0, T]} \otimes I] \tilde{U}(\tau, 0), \quad (32)$$

where $\mathbb{E}_{[0, T]} [A]$ is the expectation $\mathbb{E}[\tilde{U}(\tau, 0)^* A \tilde{U}(\tau, 0)]$.

Now if we fix arbitrary states $|\psi\rangle$ and $|\varphi\rangle$ for the probe we have that

$$\begin{aligned} \tilde{U}(T, 0) |\psi\rangle \otimes |\Omega\rangle \otimes |\varphi\rangle &= \tilde{U}(T, \tau) V U(\tau, 0) |\psi\rangle \otimes |\Omega\rangle \otimes |\varphi\rangle \\ &= \tilde{F}(T, \tau) V \tilde{F}(\tau, 0) |\psi\rangle \otimes |\Omega\rangle \otimes |\varphi\rangle \end{aligned}$$

where $\tilde{F}(t_2, t_1) = F(t_2, t_1) \otimes I$; compare (17).

Using the Theorem, we may write

$$p(m, \tau) = \frac{\text{tr} \left\{ \hat{\rho}_{\text{sys}}(\tau) \otimes \rho_{\text{probe}} V^* [E^Y(T, \tau) \otimes P_m] V \right\}}{\text{tr} \left\{ \hat{\rho}_{\text{sys}}(\tau) \otimes \rho_{\text{probe}} V^* [E^Y(T, \tau) \otimes I] V \right\}}, \quad (33)$$

where $F^Y(t_2, t_1) = \vec{T} e^{\int_{t_1}^{t_2} [\sum_k L_k dY_k(t) + K dt]}$, as in (18), $\hat{\rho}_{\text{sys}}(\tau)$ is the solution of the filter equation (13) *without* the probe, and

$$E^Y(t_2, t_1) = F^Y(t_2, t_1)^* F^Y(t_2, t_1). \quad (34)$$

Remark 9 *It is possible - and computationally preferable - to replace $\hat{\rho}_t$ in (33) with its unnormalized version ρ_t . Indeed, this is what is presented in [5].*

It is possible to use a reversed Markov process description here to compute $E^Y(T, t)$ as a function of the earliest time t . Let us introduce the notation of a backward (past-pointing) Itô increment

$$\overleftarrow{d}X(t) = X(t) - X(t - dt), \quad (35)$$

where $dt > 0$, then

$$\overleftarrow{d}F^Y(T, t) = F^Y(T, t) \left(\sum_k L_k \overleftarrow{d}Y_k(t) + K dt \right),$$

from which we find using the Itô calculus

$$\overleftarrow{d}E^Y(T, t) = \mathcal{L}\left(E^Y(T, t)\right) dt + \sum_k \left(E^Y(T, t)L_k + L_k^*E^Y(T, t)\right) \overleftarrow{d}Y_k(t). \quad (36)$$

Note that the equation (36) is structurally identical to the Belavkin-Zakai equation (19) except for the fact that the former now involves backwards Itô increments.

In their treatment, Gammelmarm et al. take the pair $(\rho_\tau, E^Y(T, \tau))$ to constitute *the* conditioned state of the system (*and* probe) at intermediate time τ given the measured quadratures over the time $[0, T]$. Their expression has the following alluring feature: ρ_τ is the unnormalized state which depends on the past measurements $\mathfrak{Y}_{[0, \tau]}$ and satisfies the Belavkin-Zakai equation (19) in τ with initial condition ρ_0 while $E^Y(T, \tau)$ is an effect which depends on the future measurements $\mathfrak{Y}_{(\tau, T]}$ and satisfies a time reversed Belavkin-Zakai (36) in τ with terminal condition $E^Y(T, T) = I$.

3 Multitime Interventions

The generalization to several intervention measurements is fairly straightforward at this stage. We consider r measurements all made by independent probes: the total probe space will then be the tensor product of the individual probe spaces. The measurements will take place at times $\tau_1 < \tau_2 < \dots < \tau_r$ in the time interval $[0, T]$ and for the k th probe we will measure an observable

$$M_k = \sum_m m P_{M_k}(m).$$

For convenience, we will drop the tensor product symbols. The various projections P_{M_k} will then by assumption commute with each other, and with the system and bath observables.

Furthermore we will assume that the entanglement of the system and the k th probe immediately before time τ_k will be implemented by a unitary V_k . We seek the probability $p(m_1, \tau_1; \dots; m_r, \tau_r)$ for a given sequence (m_1, \dots, m_r) of measurements conditional on the quadrature measurements over time $[0, T]$.

Let us introduce the mapping

$$\mathcal{G}_k[A] = V_k^* [\tilde{F}^Y(t_{k+1}, t_k)^* \otimes I] A [\tilde{F}^Y(t_{k+1}, t_k) \otimes I] V_k \quad (37)$$

for $k = 0, \dots, r$, where we understand that $t_0 = 0, t_{r+1} = T$ and $V_0 = I$.

At this stage, we should add that the conditioning results in all bath operators appearing being diagonal in terms of the measured quadratures. As such we can adopt the view that the family of commuting observables $Y_k(t)$ be simply considered as classical stochastic processes. As a result, the \mathcal{G}_k defined in (37)

are super-operators on the system-probe observables, while the V_k and the I in (38) may be understood as acting on the system-probe Hilbert space. The multitime probability, $p(m_1, \tau_1; \dots; m_r, \tau_r)$, is then

$$\frac{\text{tr}\left\{\rho_{\text{sys}} \otimes \rho_{\text{probe}} \mathcal{G}_0\left(\mathcal{G}_1\left(\dots \mathcal{G}_{r-1}\left(\mathcal{G}_r(P_r)P_{r-1}\right)\dots P_1\right)\right)\right\}}{\text{tr}\left\{\rho_{\text{sys}} \otimes \rho_{\text{probe}} \mathcal{G}_0\left(\mathcal{G}_1\left(\dots \mathcal{G}_{r-1}\left(\mathcal{G}_r(I)I\right)\dots I\right)\right)\right\}}. \quad (38)$$

4 Conclusion

The question of estimating the past is problematic in quantum theory. For instance, it does not make sense to estimate which slit a quantum particle went through in a two slit experiment if you have not measured it. We know Bayes Theorem does not work, so a blind use of the conditioning rules is unwarranted. The question as to when we may legitimately use (as opposed to misuse) estimation theory for quantum retrodiction is one that requires a detailed analysis.

The questions posed in the physics literature are of considerable interest as they push to concepts of quantum measurement and estimation to new areas. They add new layers to the traditional notions of quantum states and observables. Fortunately, the simpler intuition comes out intact for the class problems considered here, but this required a subtle analysis of what is actually going on.

In the paper, we show the consistency of the condition rule for (multiple) external measurements on an open quantum system during a period where it is undergoing simultaneous monitoring.

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References

- [1] V.P. Belavkin, Quantum filtering of Markov signals with white quantum noise. Radiotekhnika i Elektronika, **25**1445-1453 (1980); V.P. Belavkin, Quantum stochastic calculus and quantum nonlinear filtering. Journal of Multivariate Analysis, **42**, 171-201 (1992); L.M. Bouten, M. I. Guță, and H. Maassen, Stochastic Schrödinger equations. J. Phys. A, **37**, 3189-3209 (2004); L. Bouten, R. van Handel, M. R. James, An introduction to quantum filtering SIAM J. Control Optim., **46**, 2199-2241 (2007).
- [2] M. Tsang, Time-Symmetric Quantum Theory of Smoothing, Phys. Lett. **102**, 250403 (2009).
- [3] M. Tsang, Optimal waveform estimation for classical and quantum systems via time-symmetric smoothing, Phys. Rev. A **80**, 033840 (2009).

- [4] T.A. Wheatley, D.W. Berry, et al., Adaptive optical phase estimation using time-symmetric quantum smoothing, *Phys. Rev. Lett.* **104**, 93601 (2010).
- [5] S. Gammelmark, B. Julsgaard, and K. Mølmer, Past Quantum States of a Monitored System. In: *Phys. Rev. Lett.* **111**, 160401 (2013)
- [6] T.A. Wheatley, M. Tsang, I.R. Petersen, E.H. Huntington, Improved mirror position estimation using resonant quantum smoothing, *European Physics Journal Of Quantum Technology* 2, 13 (2015).
- [7] P. Campagne-Ibarcq, et al., Observing interferences between past and future quantum states in resonance fluorescence, *Phys. Rev. Lett.* **112**, 1-5 (2014).
- [8] L.P. Garcia-Pintos, J. Dressel, Probing quantumness with joint continuous measurements of noncommuting qubit observables, *PRA* **94**, 062119, December (2016); L.P. Garcia-Pintos, J. Dressel, Past observable dynamics of a continuously monitored qubit *PRA* **96**, 062110, December (2017).
- [9] S. Hacohe-Gourgy, L. Martin, E. Flurin, V.V. Ramamesh, K.B. Whaley, I. Siddiqi, Dynamics of simultaneously measured non-commuting observables, *Nature* **538**, 491-494, October (2016).
- [10] Q. Ficheux, S. Jezouin, Z. Leghtas, B. Huard, Dynamics of a qubit while simultaneously monitoring its relaxation and dephasing, *Nature Communications*, **9**, 1926, (2018).
- [11] Z.K. Mineev, et al., To catch and reverse a quantum jump mid-flight, *Nature*, **570**, pp. 200204 (2019).
- [12] H. Amini, M. Mirrahimi, P. Rouchon, On stability of continuous-time quantum filters. In: 50th IEEE Conference on Decision and European Control Conference, pp. 62426247 (2011).
- [13] P. Rouchon, Fidelity is a sub-martingale for discrete-time quantum filters. *IEEE Trans. Automat. Control* 56, 27432747 (2011).
- [14] T. Benoist, C. Pellegrini, Large Time Behavior and Convergence Rate for Quantum Filters Under Standard Non Demolition Conditions *Communications in Mathematical Physics*, 331, Issue 2, pp 703723 (2014).
- [15] L. Bouten, R. van Handel, Quantum Filtering: A Reference Probability Approach, (unpublished) arXiv:math-ph/0508006 (2005).
- [16] R.L. Hudson and K.R. Parthasarathy, Quantum Ito's formula and stochastic evolutions, *Commun. Math. Phys.* **93**, 301-323 (1984).